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## A Reinvestigation of the Crystal Structure of $\text{LiIO}_3$ \*

BY ABRAHAM ROSENZWEIG

*University of New Mexico, Albuquerque, New Mexico, U.S.A.*

AND BRUNO MOROSIN

*Sandia Laboratory, Albuquerque, New Mexico, U.S.A.*

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The crystal structure of  $\text{LiIO}_3$  has been reinvestigated with the use of single-crystal  $\text{Mo } K\alpha$  diffractometer intensity data. Morphological measurements and diffraction photographs establish  $P6_3$  to be the correct space group. The values for the lattice constants,  $a_0 = 5.1815$  and  $c_0 = 5.1709$ , were determined by a least-squares fit of back-reflection powder lines. The structure consists of lithium ions, surrounded by six oxygen atoms in a distorted octahedral arrangement, and discrete trigonal iodate groups with a  $1.817 \text{ \AA}$  iodine-oxygen bond distance and  $98.65^\circ$  O-I-O bond angle.

### Introduction

As a result of discussions concerning naturally occurring iodates (Kaplan, 1966), we reviewed the literature on the configuration of iodate groups in crystalline solids. It was noted that the structure of  $\text{LiIO}_3$  reported by Zachariasen & Barta (1931) (hereafter ZB) placed the iodine in an octahedral coordination with six equidistant oxygens, rather than in trigonal, pyramidal configuration as in other iodates. That the structure might merit reinvestigation was pointed out by Wyckoff (1960). This anomalous behavior of iodine in several alkali metal iodates has also been noted by Ibers (1956). On examination of the original paper of ZB, it became apparent that the structure was interpreted largely from powder diffraction data, and that the crystals available were extremely small. It is then not too surprising that an incorrect space group was chosen and erroneous positions were deduced for the lighter atoms. Since occasionally discrete molecules rearrange in the solid state when crystallized from gaseous or liquid state, (e.g. gaseous  $\text{Al}_2\text{Cl}_6$  and infinite molecular

sheets in crystalline  $\text{AlCl}_3$ ), the structure of  $\text{LiIO}_3$  merits reinvestigation.

Our structure results are reported here.

### Experimental

Crystals of lithium iodate were prepared in the following manner: (1) stoichiometric quantities of reagent grade iodic acid and lithium carbonate were dissolved in water; (2) the solution was then warmed to remove any remaining carbon dioxide; (3) the solution was allowed to evaporate at room temperature until crystals began to form.

The crystals obtained are brilliant pseudo-prisms up to 0.5 mm long and 0.2 mm in diameter. Careful examination reveals that the apparent prisms are actually portions of extremely steep hexagonal pyramids. Morphological measurements made with a reflecting goniometer yielded results which permit the indexing of all but the very steep faces. The polar angle of these steep faces is so great that indexing becomes ambiguous. The approximate index of these faces, orienting them so that they taper toward the negative end of  $\mathbf{c}$ , is  $\{20 \cdot 0 \cdot \bar{2}0 \cdot \bar{1}\}$ . These are terminated by the forms  $\{10\bar{1}1\}$ ,  $\{10\bar{1}\bar{1}\}$ , and very small faces of  $\{12\bar{3}\bar{2}\}$ . Mor-

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phological measurements are listed in Table 1. The morphology is illustrated in Fig. 1 and is characteristic of the point group 6 ( $C_6$ ) rather than 622 ( $D_6$ ) as reported by ZB.

Table 1. Morphological data for  $\text{LiIO}_3$

Form	$\varphi$	$\rho$
$\{10\bar{1}1\}$	$30^\circ$	$47^\circ 27'$
$\{10\bar{1}\bar{1}\}$	$30^\circ$	$-47^\circ 27'$
$\{20\ 0\ \bar{2}0\ \bar{1}\}$	$30^\circ$	$-87^\circ 20'$
$\{123\bar{2}\}$	$49^\circ 05'$	$-55^\circ 13'$

The polar character of the  $c$  axis is further established by the piezoelectric property of the crystals. Crystals which were held between two copper strips develop a positive polarity at the larger end upon compression.

The crystals of  $\text{LiIO}_3$  dissolve slowly in water, and microscopic observation shows that the ends of the crystals dissolve at different rates and in different ways. The form  $\{10\bar{1}1\}$  etches and rounds rapidly, while the form  $\{10\bar{1}\bar{1}\}$  disappears and is replaced by sharp faces of  $\{11\bar{2}\bar{1}\}$ . The appearance of the crystal after about 15 seconds in water at  $20^\circ\text{C}$  is shown in Fig. 2. Broken fragments display similar behavior.

The room-temperature unit-cell dimensions for  $\text{LiIO}_3$  were obtained by least-squares fit of 12 back-reflection powder lines obtained in a 114.5 mm Norelco powder camera with  $\text{Cr } K\alpha$  radiation. These values ( $\lambda$  for

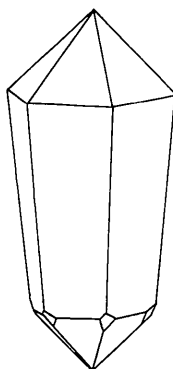


Fig. 1. Morphology of  $\text{LiIO}_3$  grown from a saturated water solution.

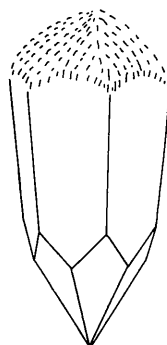


Fig. 2. The modified morphology of water-etched  $\text{LiIO}_3$  crystals.

$\text{Cr } K\alpha_1 = 2.28962 \text{ \AA}$ ),  $a_0 = 5.4815 \pm 3^*$ ,  $c_0 = 5.1709 \pm 4$ , compare favorably with previous values of  $a_0 = 5.469 \pm 3$ ,  $c_0 = 5.155 \pm 3$  (ZB) and  $a_0 = 5.481$ ,  $c_0 = 5.172$  (NBS, 1957). Systematic absence of  $000l$  reflections except for  $l = 2n$ , the morphology, and physical properties establish the space group  $P6_3$ .

From typical crystal specimens described above, a pseudo-prism, approximately 0.08 mm maximum diameter and 0.32 mm length, was selected for intensity measurements and mounted on a Pyrex rod with the  $c$  axis parallel to the axis of the rod. A Picker diffractometer equipped with a Datex semiautomatic remote control system, an Electronics and Alloys full-circle crystal orienter, a scintillation counter and a Mo X-ray tube were used for the intensity measurements. Intensities were measured at room temperature by the stationary-crystal stationary-counter technique with balanced Zr and Y filters for 257 independent reflections, of which 14 were considered unobservable and assigned a minimum value (3 times the square root of the sum transmitted by the Zr and Y filters for the background count). A small correction was added to the background (Y filter) count which was obtained by measuring the 'band width' passed (Compton scattering, thermal diffuse scattering and parasitic scattering) by the filter pair for a particular  $2\theta$  value (Burbank, 1965). The data were corrected for Lorentz and polarization factors. The absorption coefficient for Mo  $K\alpha$  is  $117 \text{ cm}^{-1}$ ; hence absorption corrections are quite small. However, cylindrical absorption corrections approximating the pseudo-prism were applied; the values of the absorption factors ranged from 2.07 to 2.25. Structure factors were calculated with the use of the lithium, oxygen and iodine scattering factors given in *International Tables for X-ray Crystallography* (1962, Tables 3.3.1A and 3.3.1B). Iodine dispersion corrections ( $\Delta f' = -0.6$ ;  $\Delta f'' = 2.2$ ) were included by the method of Templeton (1955) in all later stages of refinement. Calculations were carried out on CDC 1604 and 3600 computers.

### Structure determination

The distribution of intensities ( $h h \bar{2} h l$  with  $l = 2n + 1$  very weak and  $h0hl$  all about equal except when  $h = 3$ ) suggests that iodine atoms occupy ( $b$ ) sites in space group  $P6_3$ . These positions are equivalent to those reported by ZB. A three-dimensional Fourier synthesis was calculated with the use of phases determined by the iodine atoms located in ( $b$ ) sites with  $z$  arbitrarily set at zero in all calculations. This Fourier synthesis indicated locations for the oxygen atoms as well as a set of centrosymmetrically related 'ghosts'. Refinement of the three oxygen positional and two isotropic thermal parameters as well as the one scale factor by least-squares methods [unit weights were assigned all reflections except for unobserved reflections

\* Throughout this paper, the value of the error corresponds to the last significant digits in the function value.



a shortening of the shorter Li–O separations given above. Refinement of model *A* results in a lithium parameter of *z* equal to 0.1286 (with Li–O values of 1.71 and 2.78 Å); fractional positional parameters for oxygen differed by  $1 \times 10^{-5}$  and  $B_{ij}$  (oxygen) differed by a maximum of 0.1 $\sigma$  from values for model *B* listed in Table 2. However, one may rule out model *A* by consideration of the sum of ionic radii (using 0.60 Å for Li<sup>+</sup> and 1.32 to 1.40 Å for O<sup>2-</sup>; Pauling, 1960).

The final positional and thermal parameters (Table 2) result in a discrepancy index of 0.049. Anisotropic thermal parameters,  $B_{ijn}$ , are of the form  $\exp(-\frac{1}{4} \sum_{i=1}^3 \sum_{j=1}^3 B_{ijn} h_i h_j a_i^* a_j^*)$ . Table 3 lists the observed and calculated structure factors scaled by 10 times as indicated. Table 4 lists the interatomic separations and bond angles. Other reported I–O bond lengths and O–I–O angles have been summarized by Ibers (1956), and range from acceptable values of 1.77 to 1.86 (average 1.82) found in NaIO<sub>3</sub> (van Eck and MacGillavry, 1943), Ce(IO<sub>3</sub>)<sub>4</sub> (Cromer & Larson, 1956), Ce(IO<sub>3</sub>)<sub>4</sub> · H<sub>2</sub>O (Ibers, 1956), and HIO<sub>3</sub> (Rogers & Helmholz,

1941) to improbable values of 2.08 Å (Naray-Szabo & Neugebauer, 1947) or larger (suggested by octahedral iodine in the iodates of cesium, rubidium and ammonium ion).

It is unfortunate in a crystal whose space group is unambiguously established that the iodine atoms are related to each other by a point of inversion, and hence one is not able to utilize their dispersion to resolve the absolute configuration of the packing.

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Table 4. Interatomic separations in LiIO<sub>3</sub>

	Separations	Standard deviation
Li–O	2.038, 2.220 Å	0.06 Å
I–O (bonded)	1.817	0.017
I–O (next neighbor)	2.873	0.017
O–O (intra IO <sub>3</sub> <sup>-</sup> groups)	2.756	0.024
O–O (inter IO <sub>3</sub> <sup>-</sup> groups)	2.925, 3.085	0.024
O–I–O	98.65°	0.6°

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## Cation Distribution and Oxygen Parameter in Magnesium Gallate, MgGa<sub>2</sub>O<sub>4</sub>

BY JAMES E. WEIDENBORNER, N. R. STEMPLE AND Y. OKAYA

*IBM Watson Research Center, Yorktown Heights, New York, U.S.A.*

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Magnesium gallate, MgGa<sub>2</sub>O<sub>4</sub>, a partially inverted spinel, belongs to space group *Fd3m* with  $a_0 = 8.286 \pm 0.003$  Å. The distribution of the di- and trivalent cations on two possible sites, 8(*a*) and 16(*d*), has been determined by the least-squares method. The data used in the refinement were obtained from single crystals by film and counter techniques; the counter data were collected manually and by CXXD, a computer-controlled X-ray diffractometer.

The magnesium ions are distributed between the two sites as follows:  $16 \pm 2\%$  in the tetrahedral *A* site and  $81 \pm 1\%$  in the remaining octahedral *B* site. The oxygen parameter as refined in the present study is  $0.3824 \pm 0.0003$ , indicating the usual shift for the inverted spinel from the ideal value of  $3/8$ .

### Introduction

Magnesium gallate, MgGa<sub>2</sub>O<sub>4</sub>, is a partially inverted spinel with space group *Fd3m*. The metal ions fall on sets of special positions 8(*a*) and 16(*d*), frequently referred to as the *A* and *B* sites respectively. The oxygen

ions occupy a set of special positions 32(*e*) with one positional parameter to be determined.

Interest in this compound developed through electron spin resonance studies of Cu<sup>2+</sup> conducted at these Laboratories (Overmeyer, 1963) using magnesium gallate as a host lattice. A substantial broadening is ob-